

Lecture 4

Boundary Conditions

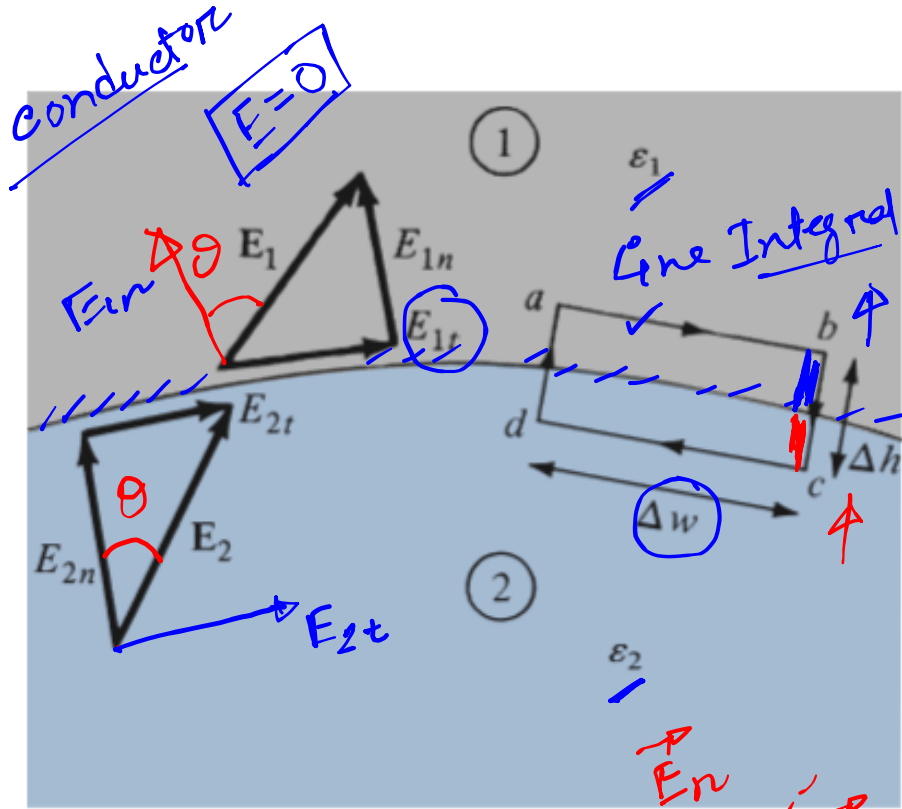
Nazmul Haque Turja

Research and Development Assistant, BUET

$$E_{1n} \neq E_{2n}$$

Electric Boundary Conditions(Dielectric-Dielectric)

$$\oint_L \mathbf{E} \cdot d\mathbf{l} = 0$$



$$0 = E_{1t} \Delta w - \cancel{E_{1n} \frac{\Delta h}{2}} - \cancel{E_{2n} \frac{\Delta h}{2}} - E_{2t} \Delta w + \cancel{E_{2n} \frac{\Delta h}{2}} + \cancel{E_{1n} \frac{\Delta h}{2}}$$

where $E_t = |\mathbf{E}_t|$ and $E_n = |\mathbf{E}_n|$. The $\frac{\Delta h}{2}$ terms cancel, and eq. (5.56) becomes

$$0 = (E_{1t} - E_{2t}) \Delta w$$

or

$$\vec{E} = \vec{E}_{tan} + \vec{E}_n$$

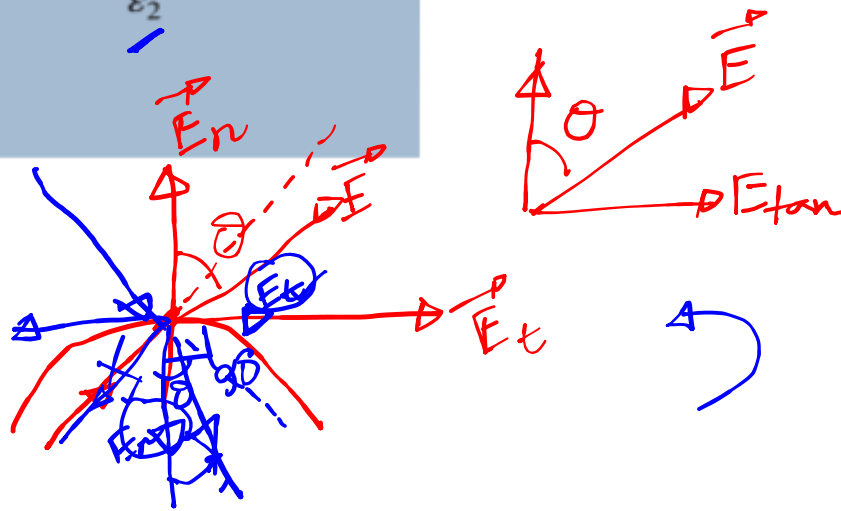
$$\mathbf{D} = \epsilon \mathbf{E}$$

$$E_{1t} = E_{2t}$$

$$\frac{D_{1t}}{\epsilon_1} = E_{1t} = E_{2t} = \frac{D_{2t}}{\epsilon_2}$$

$$\frac{D_{1t}}{\epsilon_1} = \frac{D_{2t}}{\epsilon_2}$$

$\theta = \text{Normal} \wedge \text{Vector}$





Electric Boundary Conditions (Dielectric-Dielectric)

$$\oint_S \mathbf{D} \cdot d\mathbf{S} = Q_{\text{enc}}$$

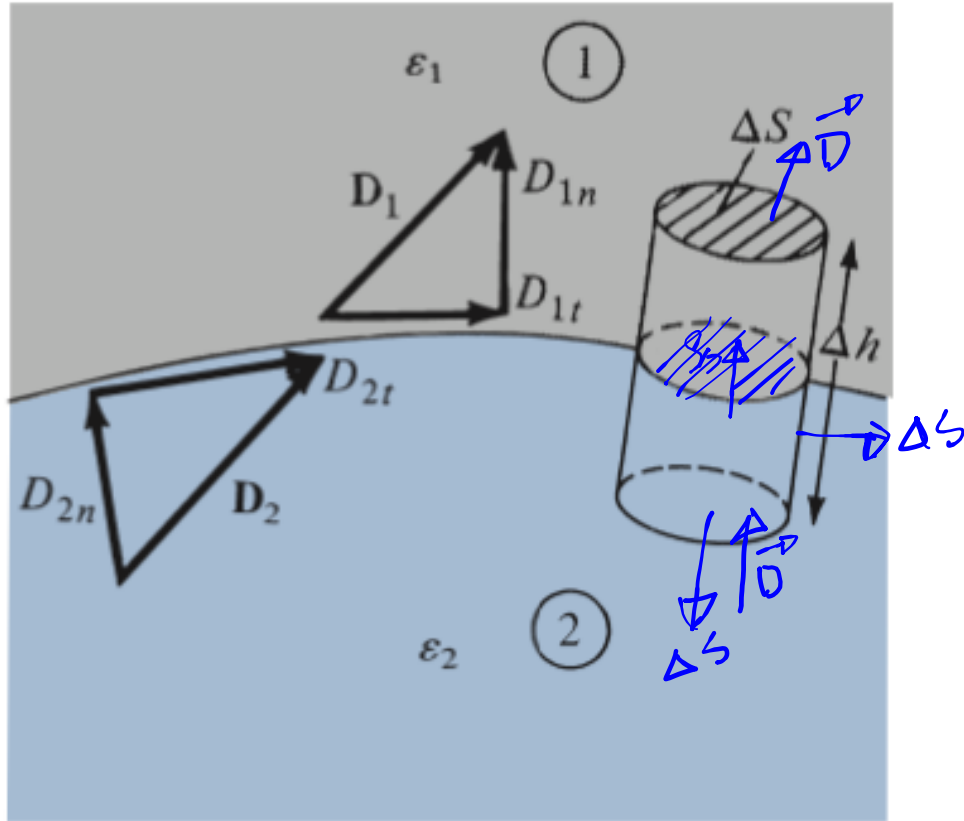
$$\rho_s = \frac{Q_{\text{enc}}}{\Delta S}$$

$$Q = \rho_s \Delta S$$

$$\Delta Q = \rho_s \Delta S = D_{1n} \Delta S - D_{2n} \Delta S$$

$$\cos 90^\circ = 0$$

$$D_{1n} - D_{2n} = \rho_s$$



(b)

where ρ_s is the free charge density placed deliberately at the boundary. It should be borne in mind that eq. (5.59) is based on the assumption that \mathbf{D} is directed from region 2 to region 1 and eq. (5.59) must be applied accordingly. If no free charges exist at the interface (i.e., charges are not deliberately placed there), $\rho_s = 0$ and eq. (5.59) becomes

$$D_{1n} = D_{2n} \quad (5.60)$$

Thus the normal component of \mathbf{D} is continuous across the interface; that is, D_n undergoes no change at the boundary. Since $\mathbf{D} = \epsilon \mathbf{E}$, eq. (5.60) can be written as

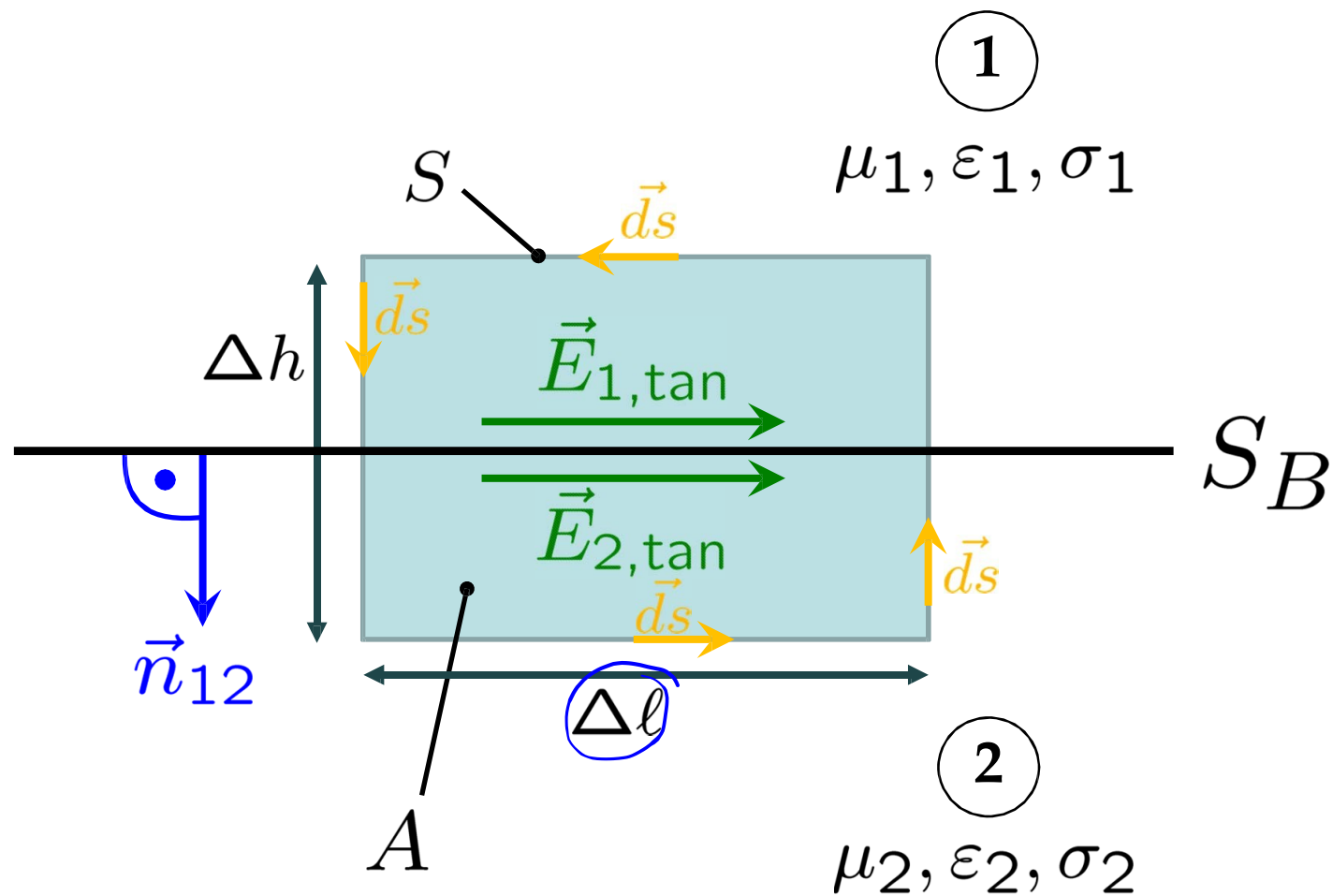
$$\epsilon_1 E_{1n} = \epsilon_2 E_{2n} \quad (5.61)$$

$$\Rightarrow E_{1n} = \epsilon_2 / \epsilon_1 E_{2n}$$

Boundary Conditions

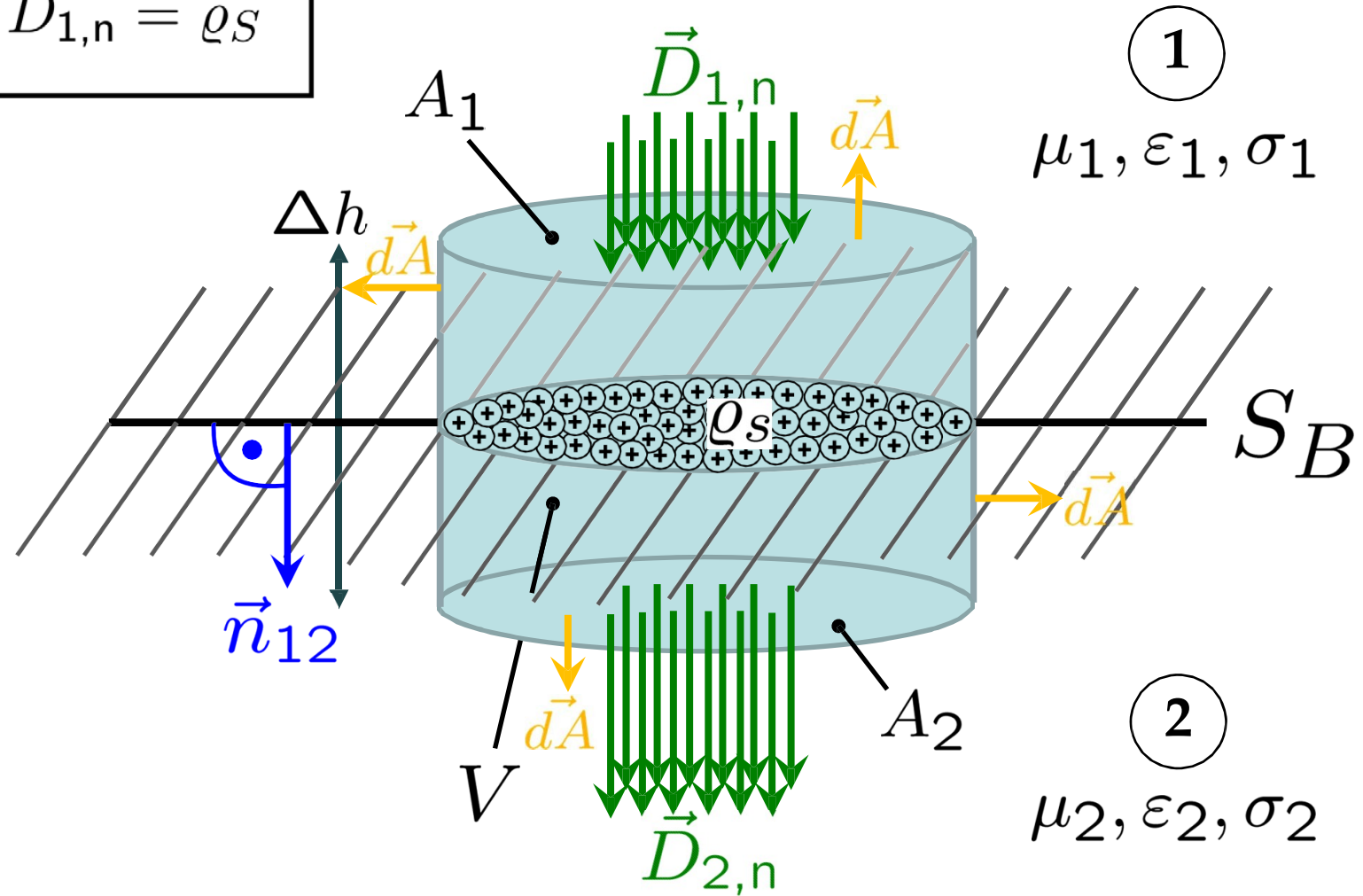
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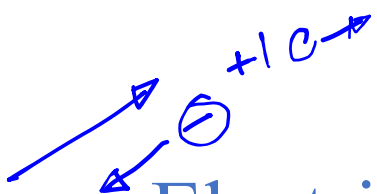
$$\vec{E}_{2,\text{tan}} = \vec{E}_{1,\text{tan}}$$



Boundary Conditions

$$D_{2,n} - D_{1,n} = \rho_S$$

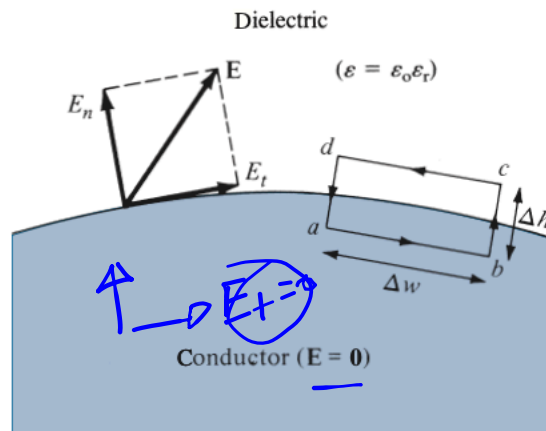
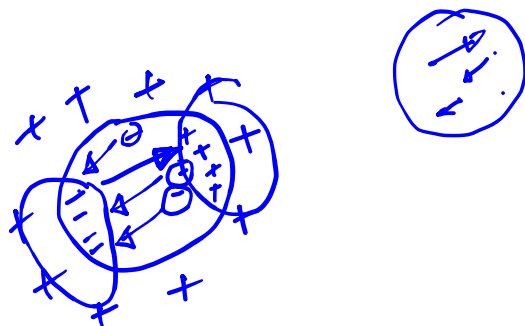




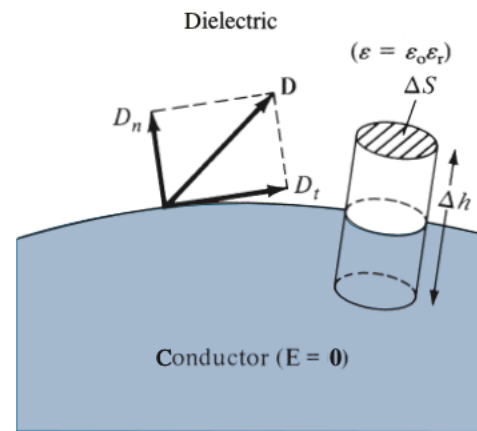
Electric Boundary Conditions(Conductor-Dielectric)

$$\oint_L \mathbf{E} \cdot d\mathbf{l} = 0$$

$$\oint_S \mathbf{D} \cdot d\mathbf{S} = Q_{\text{enc}}$$



(a)



(b)

$$D_2 \Delta S - D_1 \Delta S = 0$$

abceda

$$0 = 0 \cdot \Delta w + 0 \cdot \frac{\Delta h}{2} + E_n \cdot \frac{\Delta h}{2} - E_t \cdot \Delta w - E_n \cdot \frac{\Delta h}{2} - 0 \cdot \frac{\Delta h}{2}$$

As $\Delta h \rightarrow 0$,

$$E_t = 0$$

$$E_{t1} = E_{t2} = 0$$

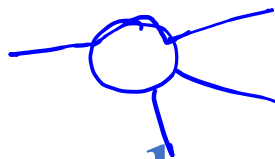
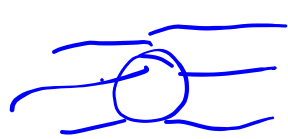
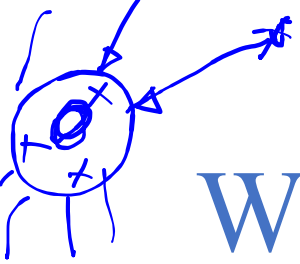
$$\Delta Q = D_n \cdot \Delta S - 0 \cdot \Delta S$$

because $\mathbf{D} = \epsilon \mathbf{E} = \mathbf{0}$ inside the conductor. Equation (5.68) may be written as

$$D_n = \frac{\Delta Q}{\Delta S} = \rho_s$$

or

$$D_n = \rho_s$$



What is a perfect conductor?

$$R \propto L$$
$$R \propto \frac{1}{A}$$

$$R = \frac{\rho L}{A}$$

Thus under static conditions, the following conclusions can be made about a perfect conductor:

1. No electric field may exist *within* a conductor; that is, considering our conclusion in Section 5.4,

$$\rho = \frac{L}{A} = R$$

(resistivity)

$$E = -\nabla V$$

$$\rho_v = 0, \quad E = 0$$

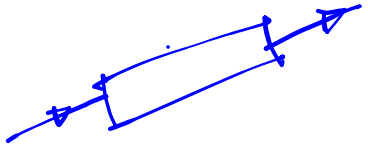
(5.70)

$$\frac{10V}{10V}$$

2. Since $E = -\nabla V = 0$, there can be no potential difference between any two points in the conductor; that is, a conductor is an equipotential body.
3. An electric field E must be external to the conductor and must be normal to its surface; that is,

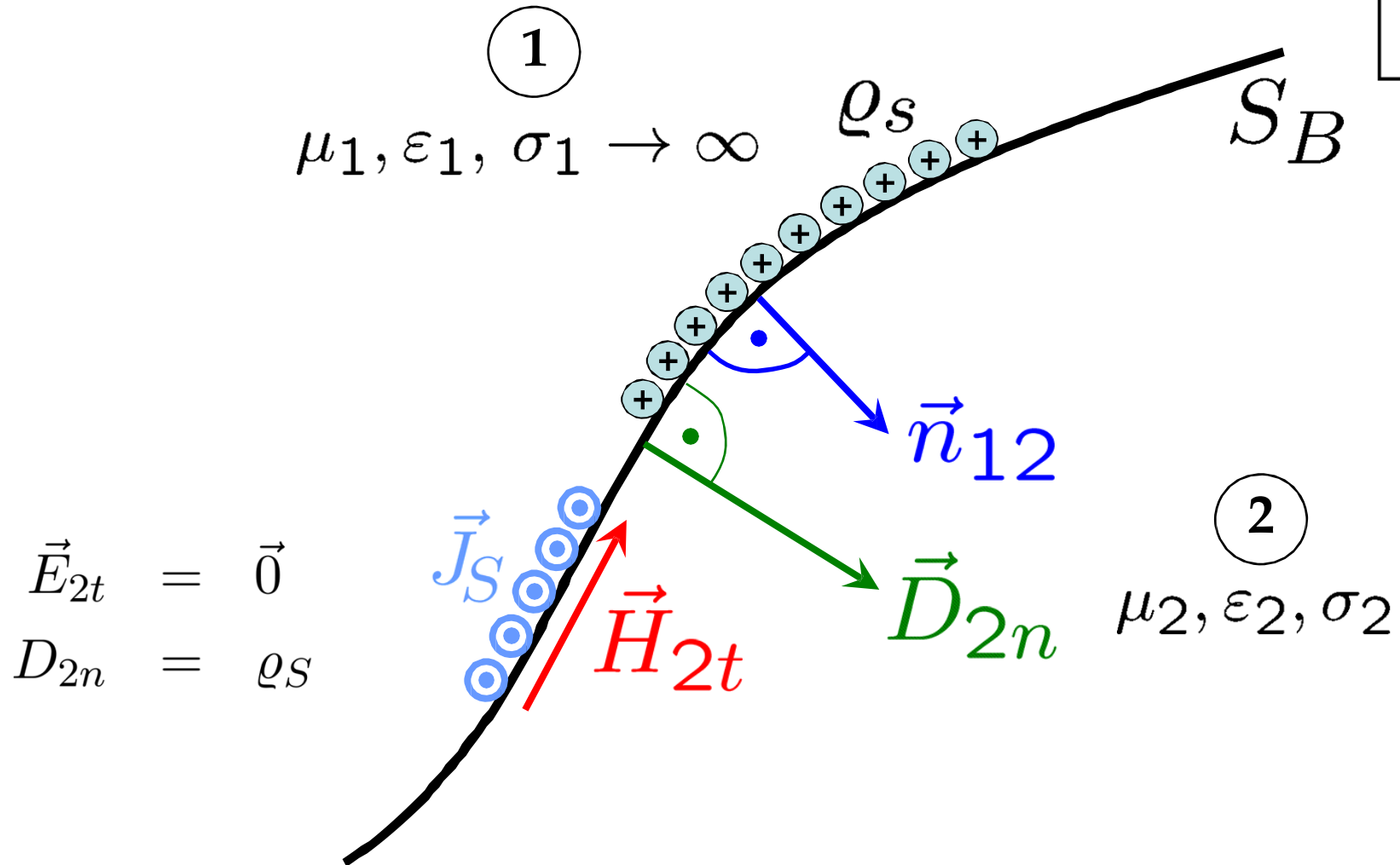
$$D_t = \epsilon_0 \epsilon_r E_t = 0, \quad D_n = \epsilon_0 \epsilon_r E_n = \rho_s$$

(5.71)



Maxwell's Equations

Boundary conditions at a perfect electric conductor:



$$\vec{E}_{2,\text{tan}} = \vec{E}_{1,\text{tan}}$$

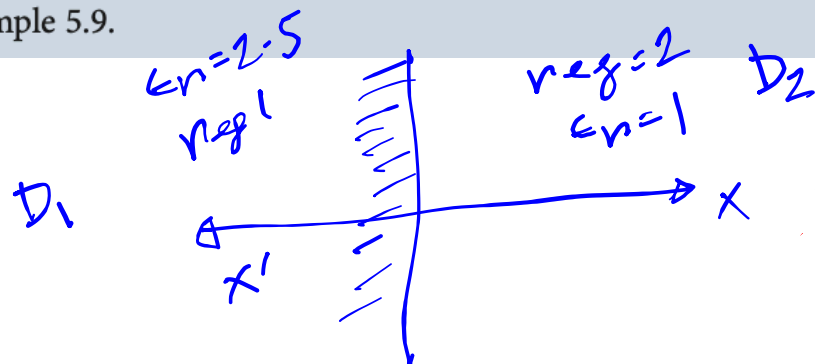
$$D_{2,n} - D_{1,n} = \rho_s$$

PRACTICE EXERCISE 5.9

A homogeneous dielectric ($\epsilon_r = 2.5$) fills region 1 ($x < 0$) while region 2 ($x > 0$) is free space.

(a) If $\mathbf{D}_1 = 12\mathbf{a}_x - 10\mathbf{a}_y + 4\mathbf{a}_z$ nC/m², find \mathbf{D}_2 and θ_2 .

(b) If $E_2 = 12$ V/m and $\theta_2 = 60^\circ$, find E_1 and θ_1 . Take θ_1 and θ_2 as defined in Example 5.9.



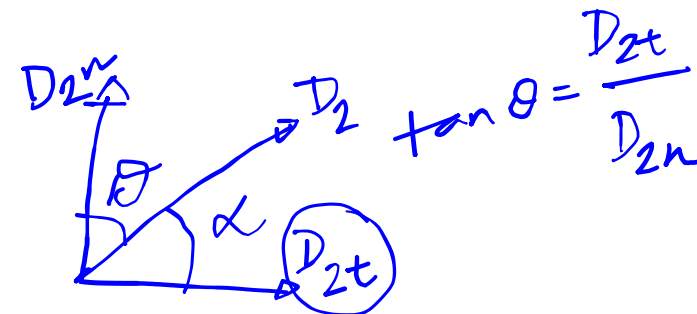
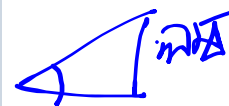
$$\tan \theta_2 = \frac{D_{2t}}{D_{2n}} = \frac{\sqrt{(-4)^2 + (1.6)^2}}{12} = 0.359 \longrightarrow \theta_2 = \underline{\underline{19.75^\circ}}$$

P. E. 5.9 (a) Since $a_n = a_x$,

$$\vec{D}_{1n} = 12\hat{a}_x, \quad D_{1t} = -10a_y + 4a_z, \quad \textcircled{D_{2n}} = D_{1n} = 12a_x$$

$$\underline{\underline{E_{2t} = E_{1t}}} \longrightarrow D_{2t} = \frac{\epsilon_2 D_{1t}}{\epsilon_1} = \frac{1}{2.5}(-10a_y + 4a_z) = -4a_y + 1.6a_z$$

$$\underline{\underline{D_2 = D_{2n} + D_{2t} = -12a_x - 4a_y + 1.6a_z \text{ nC/m}^2.}}$$



tangential angle, $\alpha = 90^\circ - \theta$

PRACTICE EXERCISE 5.9

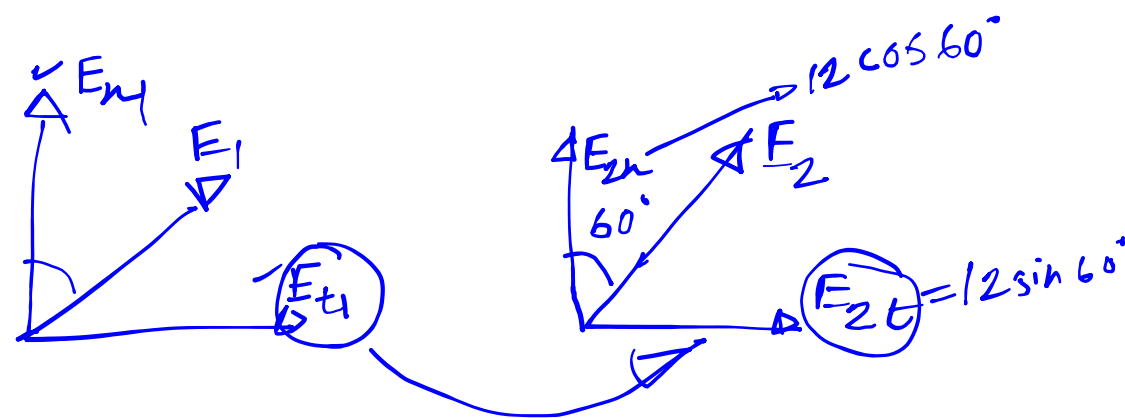
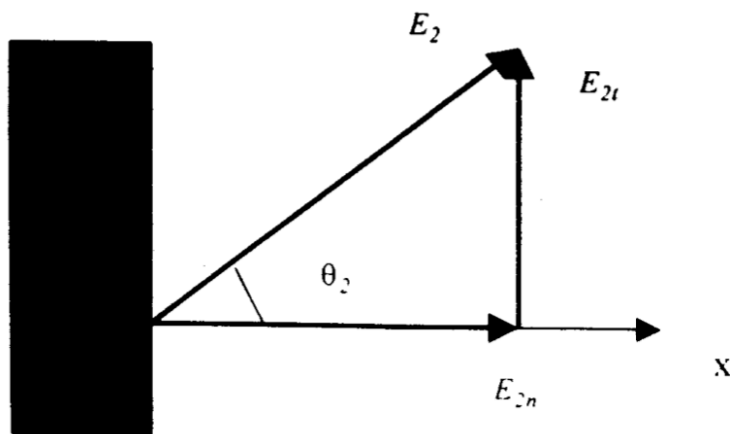
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 (b) If $E_2 = 12$ V/m and $\theta_2 = 60^\circ$, find E_1 and θ_1 . Take θ_1 and θ_2 as defined in Example 5.9.

$$\textcircled{E_2} = \dots$$

E_1

$\textcircled{b} E_{1t} = E_{2t} = E_2 \sin \theta_2 = 12 \sin 60^\circ = 10.392$



$$D_{1n} = D_{2n}$$

$$\epsilon E_{1n} = \epsilon_0 E_{2n}$$

$$E_{1n} = \frac{\epsilon_{r2}}{\epsilon_{r1}} E_{2n} = \frac{1}{2.5} 12 \cos 60^\circ = 2.4$$

$$E_1 = \sqrt{E_{1t}^2 + E_{1n}^2} = \underline{\underline{10.67}}$$

$$\tan \theta_1 = \frac{\epsilon_{r1}}{\epsilon_{r2}} \tan \theta_2 = \frac{2.5}{1} \tan 60^\circ = 4.33 \quad \longrightarrow \quad \underline{\underline{\theta_1 = 77^\circ}}$$

Note that $\theta_1 > \theta_2$.

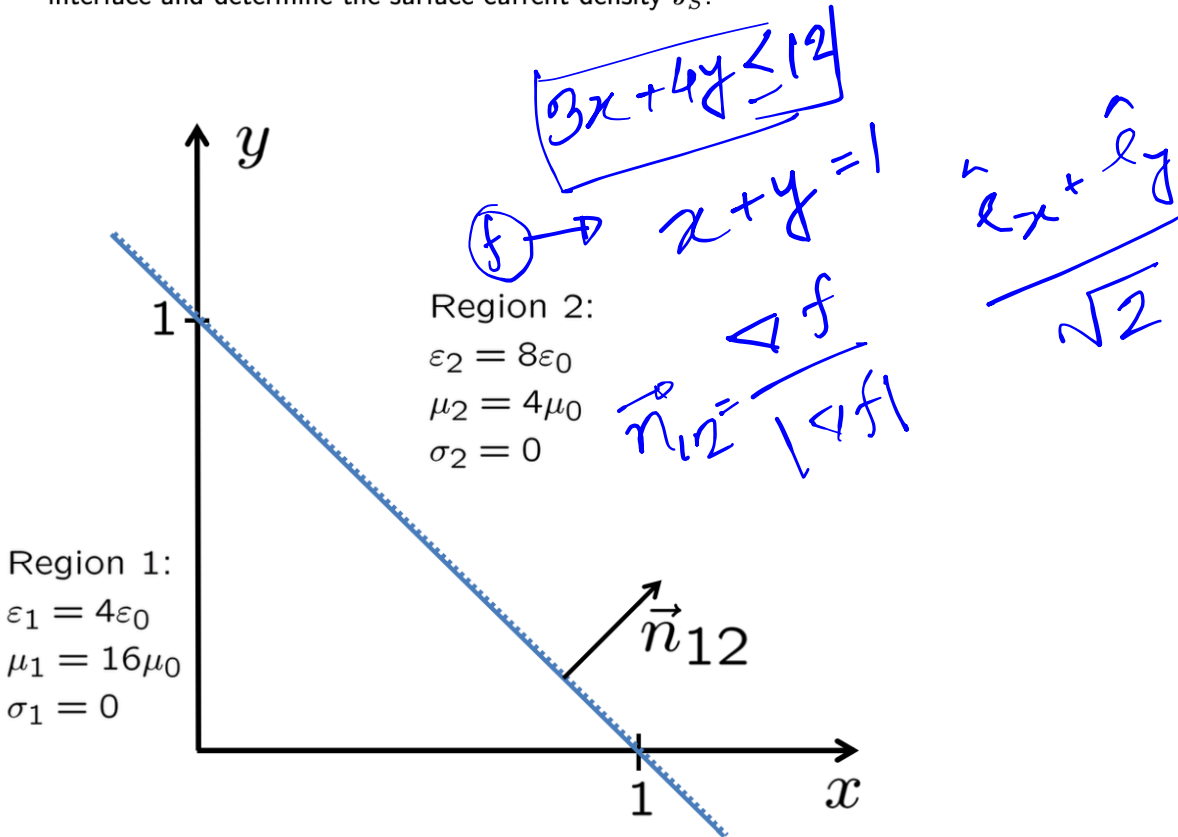
3.3 Problem 3

The three dimensional space is split into two regions filled with different materials (see Fig. 3.3.2). The electric vector field \vec{E}_1 and the magnetic vector field \vec{H}_1 , in region 1 (see Fig. 3.3.2), are given as

$\vec{E}_1 = (2.0\vec{e}_y + 3.0\vec{e}_z) \text{ V/m}$ $\vec{H}_1 = (0.1\vec{e}_x + 0.2\vec{e}_z) \text{ A/m}$

For sub-problems a) and b) there are no surface charges ρ_s on the interface.

- a) Find the electric vector field \vec{E}_2 in region 2 so that the boundary conditions at the interface are fulfilled.
- b) Find the magnetic induction \vec{B}_2 in region 2 so that the boundary conditions at the interface are fulfilled.
- c) Now the medium in region 2 is replaced by an ideal conductor with conductivity of $\sigma_2 \rightarrow \infty$ while all other parameters remain the same. Determine the surface charge density ρ_s on the interface and determine the surface current density \vec{J}_s .



a) Normal vector: $\vec{n}_{12} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix}$ ✓
 $\vec{E}_1 = \begin{pmatrix} 0 \\ 2 \\ 3 \end{pmatrix} \frac{\text{V}}{\text{m}}$; $\vec{H}_1 = \begin{pmatrix} 0.1 \\ 0 \\ 0.2 \end{pmatrix} \frac{\text{A}}{\text{m}}$

$\vec{E}_{2tan} = \vec{E}_{1tan}$
Please note: $\vec{n}_{12} \cdot \vec{E}_1 = |\vec{E}_{1n}|$; $\vec{E}_{1n} = (\vec{n}_{12} \cdot \vec{E}_1) \vec{n}_{12}$
But: $\vec{n}_{12} \times \vec{E}_1 \neq \vec{E}_{1tan}$
Instead: $\vec{E}_{1tan} = \vec{n}_{12} \times \vec{E}_1 \times \vec{n}_{12}$
Here: $\vec{n}_{12} \times \vec{E}_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \times \begin{pmatrix} 0 \\ 2 \\ 3 \end{pmatrix} \frac{\text{V}}{\text{m}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 3 \\ -3 \\ 2 \end{pmatrix} \frac{\text{V}}{\text{m}}$
 $\vec{E}_{1tan} = \vec{n}_{12} \times \vec{E}_1 \times \vec{n}_{12} = \frac{1}{\sqrt{2}} \begin{pmatrix} 3 \\ -3 \\ 2 \end{pmatrix} \times \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{\text{V}}{\text{m}} = \frac{1}{2} \begin{pmatrix} -2 \\ 2 \\ 6 \end{pmatrix} \frac{\text{V}}{\text{m}} = \begin{pmatrix} -1 \\ 1 \\ 3 \end{pmatrix} \frac{\text{V}}{\text{m}}$

$\Rightarrow \vec{E}_{2tan} = \begin{pmatrix} -1 \\ 1 \\ 3 \end{pmatrix} \frac{\text{V}}{\text{m}}$ ✓
 $|\vec{E}_{1n}| = \vec{n}_{12} \cdot \vec{E}_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 2 \\ 3 \end{pmatrix} \frac{\text{V}}{\text{m}} = \frac{2}{\sqrt{2}} \frac{\text{V}}{\text{m}} = \sqrt{2} \frac{\text{V}}{\text{m}}$
 $\vec{E}_{1n} = (\vec{n}_{12} \cdot \vec{E}_1) \vec{n}_{12} = \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{\text{V}}{\text{m}}$ ✓

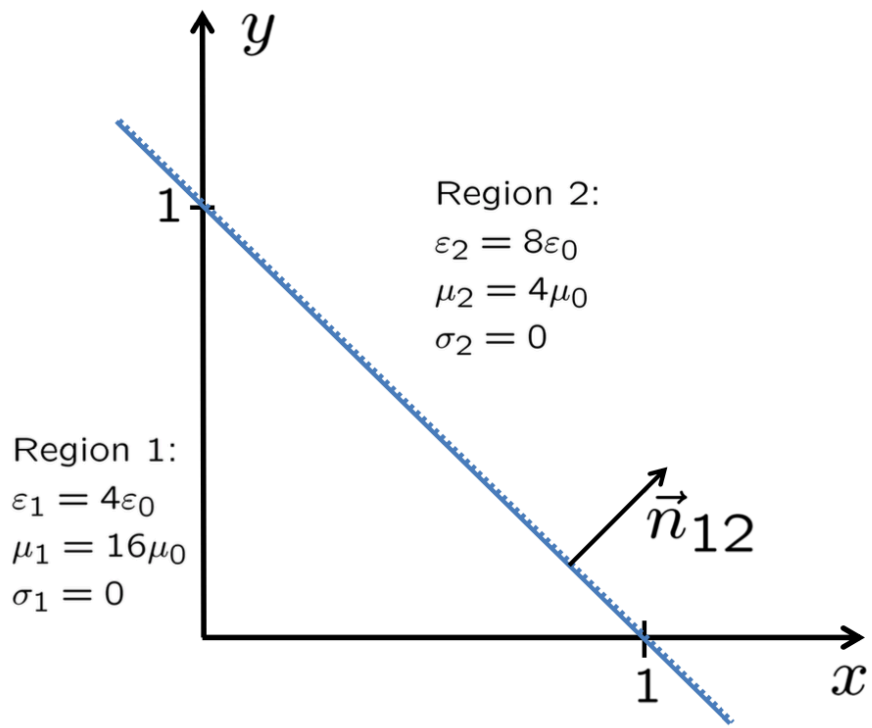
3.3 Problem 3

The three dimensional space is split into two regions filled with different materials (see Fig. 3.3.2). The electric vector field \vec{E}_1 and the magnetic vector field \vec{H}_1 , in region 1 (see Fig. 3.3.2), are given as

$$\vec{E}_1 = (2.0\vec{e}_y + 3.0\vec{e}_z) \text{ V/m} \qquad \vec{H}_1 = (0.1\vec{e}_x + 0.2\vec{e}_z) \text{ A/m}$$

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$\Rightarrow \vec{E}_{2tan} = \begin{pmatrix} -1 \\ 1 \\ 3 \end{pmatrix} \frac{V}{m}$ (10)

$|\vec{E}_{1n}| = \vec{n}_{12} \cdot \vec{E}_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 2 \\ 3 \end{pmatrix} \frac{V}{m} = \frac{2}{\sqrt{2}} \frac{V}{m} = \sqrt{2} \frac{V}{m}$

$\vec{E}_{1n} = (\vec{n}_{12} \cdot \vec{E}_1) \vec{n}_{12} = \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{V}{m}$

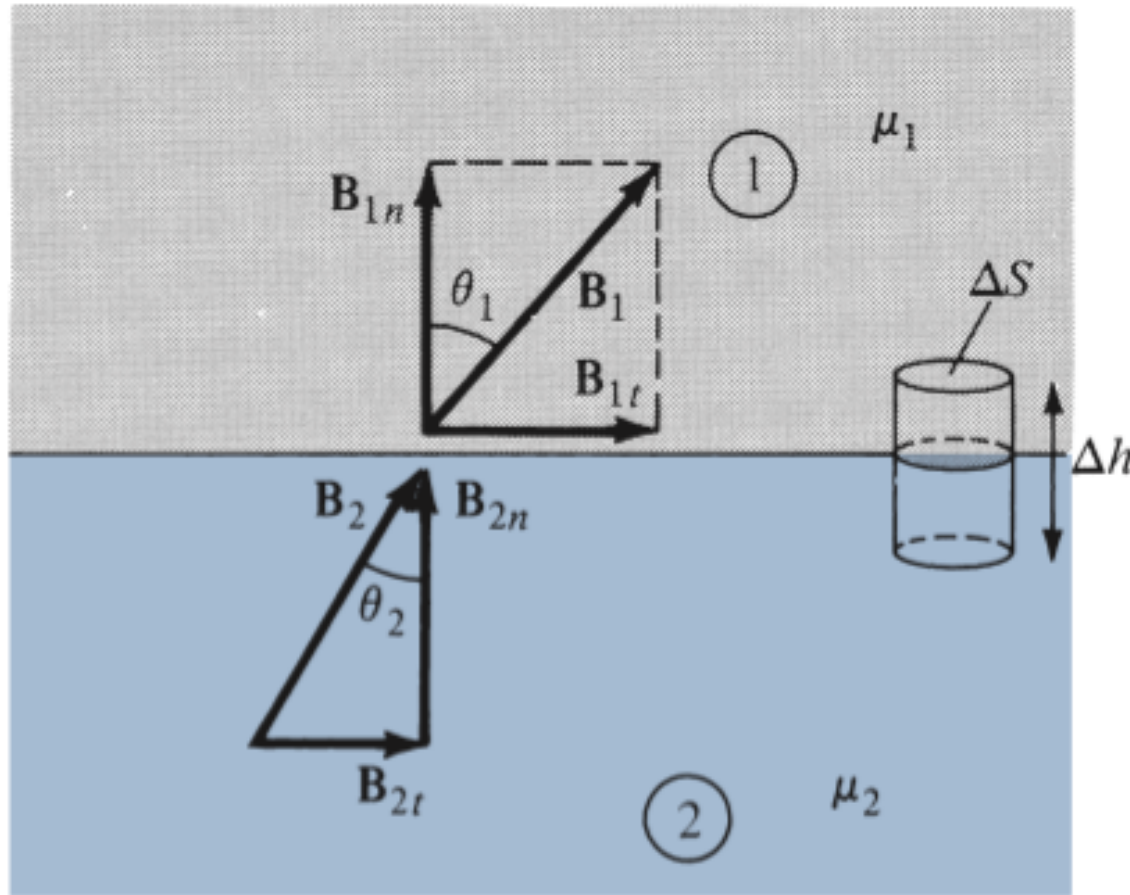
$D_{2n} = D_{1n} \Rightarrow \epsilon_2 E_{2n} = \epsilon_1 E_{1n} \Rightarrow E_{2n} = \frac{\epsilon_1}{\epsilon_2} E_{1n} = \frac{1}{2} E_{1n}$

$\Rightarrow \vec{E}_{2n} = \frac{1}{2} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{V}{m}$

$\vec{E}_2 = \vec{E}_{2tan} + \vec{E}_{2n} = \begin{pmatrix} -1 \\ 1 \\ 3 \end{pmatrix} \frac{V}{m} + \begin{pmatrix} 1/2 \\ 1/2 \\ 0 \end{pmatrix} \frac{V}{m} = \begin{pmatrix} -1/2 \\ 3/2 \\ 3 \end{pmatrix} \frac{V}{m}$

Magnetic Boundary Conditions

$$\oint_S \mathbf{B} \cdot d\mathbf{S} = \underline{0}$$



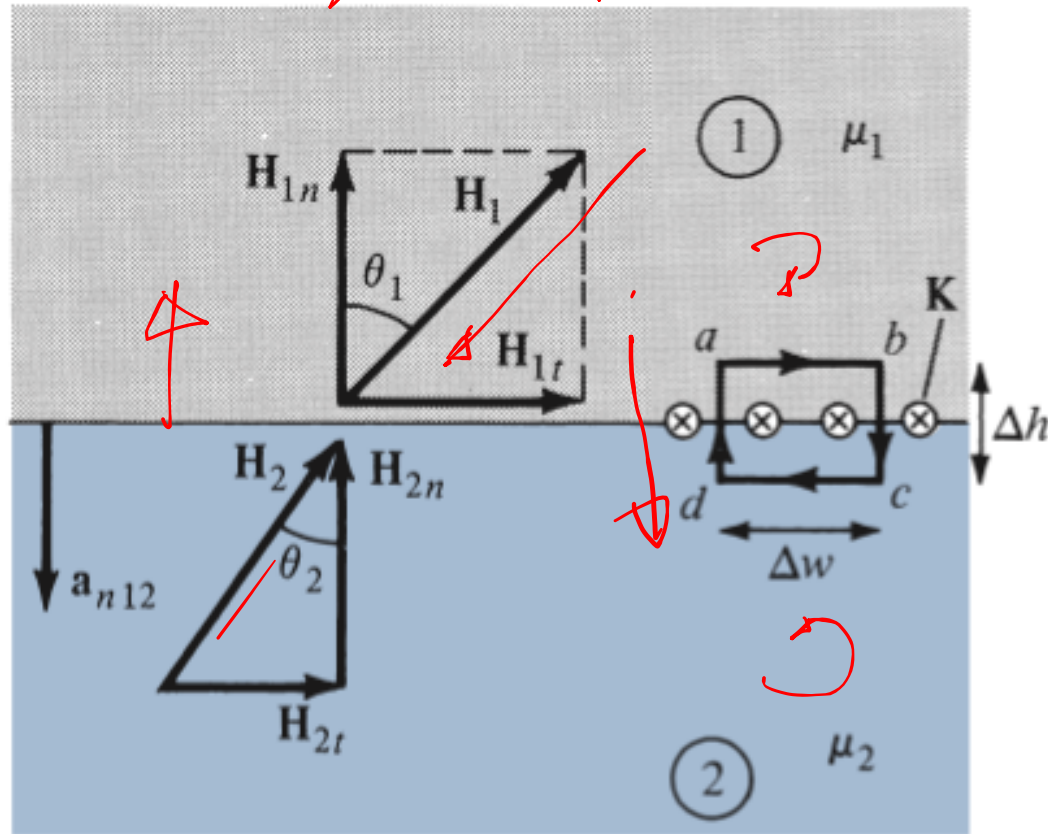
(a)

$$B_{1n} \Delta S - B_{2n} \Delta S = 0$$

\checkmark $\boxed{B_{1n} = B_{2n}}$ or $\mu_1 \mathbf{H}_{1n} = \mu_2 \mathbf{H}_{2n}$ \checkmark

Magnetic Boundary Conditions

$$H_{1t} \times \Delta w - H_{1n} \times \frac{\Delta h}{2} - H_{2n} \times \frac{\Delta h}{2} - H_{2t} \times \Delta w + H_{2n} \times \frac{\Delta h}{2} + H_{2t} \times \frac{\Delta h}{2}$$



(b)

$$\vec{H}_{1t} - \vec{H}_{2t} = \vec{K}$$

$$\oint_L \mathbf{H} \cdot d\mathbf{l} = I$$

Similarly, we apply eq. (8.39) to the closed path $abcd$ of Figure 8.16(b), where surface current \mathbf{K} on the boundary is assumed normal to the path. We obtain

$$K \cdot \Delta w = H_{1t} \cdot \Delta w + H_{1n} \cdot \frac{\Delta h}{2} + H_{2n} \cdot \frac{\Delta h}{2}$$

As $\Delta h \rightarrow 0$, eq. (8.42) leads to

$$H_{1t} - H_{2t} = K \quad (8.43)$$

This shows that the tangential component of H is also discontinuous. Equation (8.43) may be written in terms of B as

$$\frac{B_{1t}}{\mu_1} - \frac{B_{2t}}{\mu_2} = K \quad (8.44)$$

In the general case, eq. (8.43) becomes

$$(\mathbf{H}_1 - \mathbf{H}_2) \times \mathbf{a}_{n12} = \mathbf{K} \quad (8.45)$$

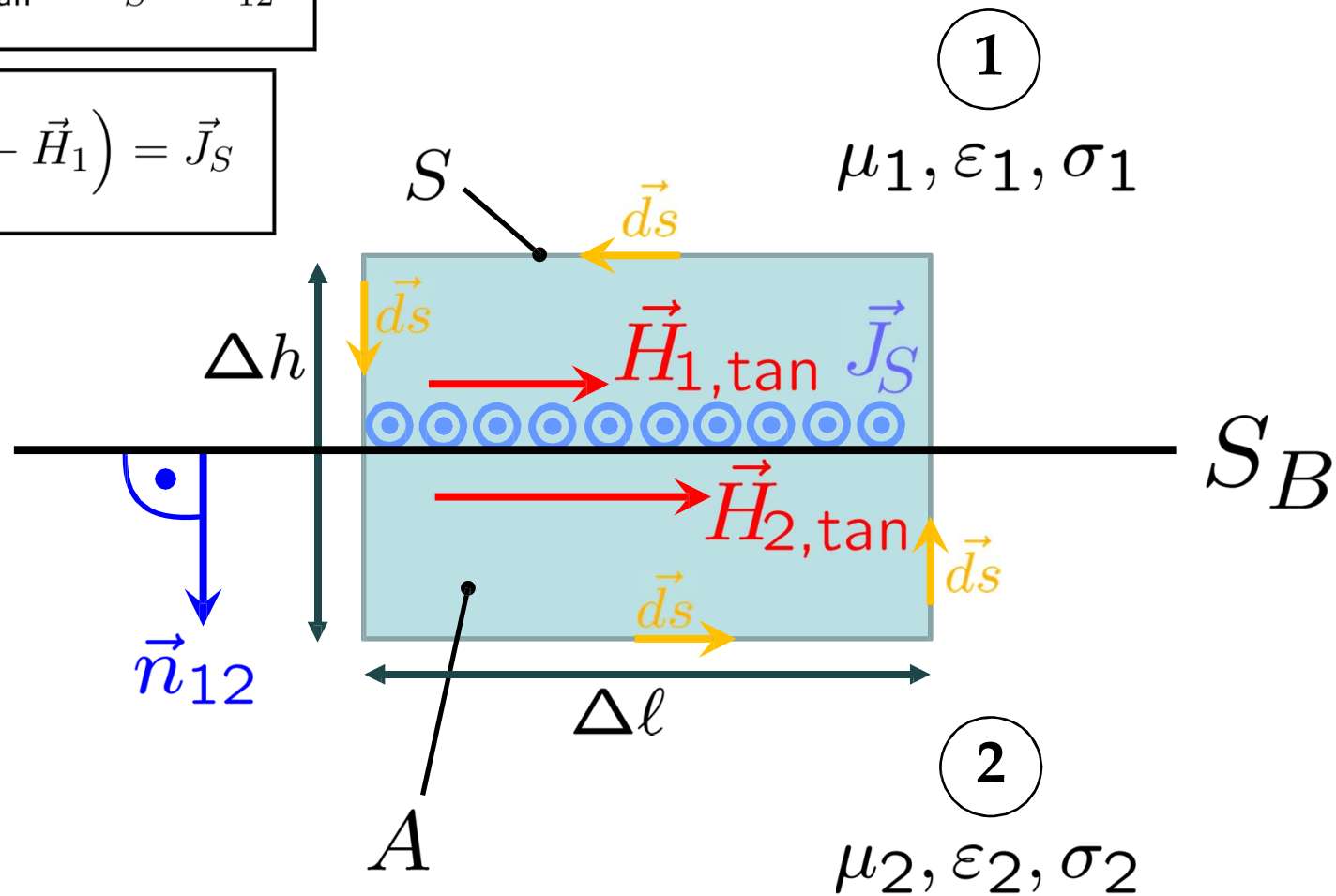
where \mathbf{a}_{n12} is a unit vector normal to the interface and is directed from medium 1 to medium 2. If the boundary is free of current or the media are not conductors (for \mathbf{K} is free current density), $\mathbf{K} = \mathbf{0}$ and eq. (8.43) becomes

$$\mathbf{H}_{1t} = \mathbf{H}_{2t} \quad \text{or} \quad \frac{B_{1t}}{\mu_1} = \frac{B_{2t}}{\mu_2} \quad (8.46)$$

Boundary Conditions

$$\vec{H}_{2,\text{tan}} - \vec{H}_{1,\text{tan}} = \vec{J}_S \times \vec{n}_{12}$$

$$\vec{n}_{12} \times (\vec{H}_2 - \vec{H}_1) = \vec{J}_S$$



Boundary Conditions

$$\rightarrow \vec{H}_{2,\text{tan}} - \vec{H}_{1,\text{tan}} = \vec{J}_S \times \vec{n}_{12}$$

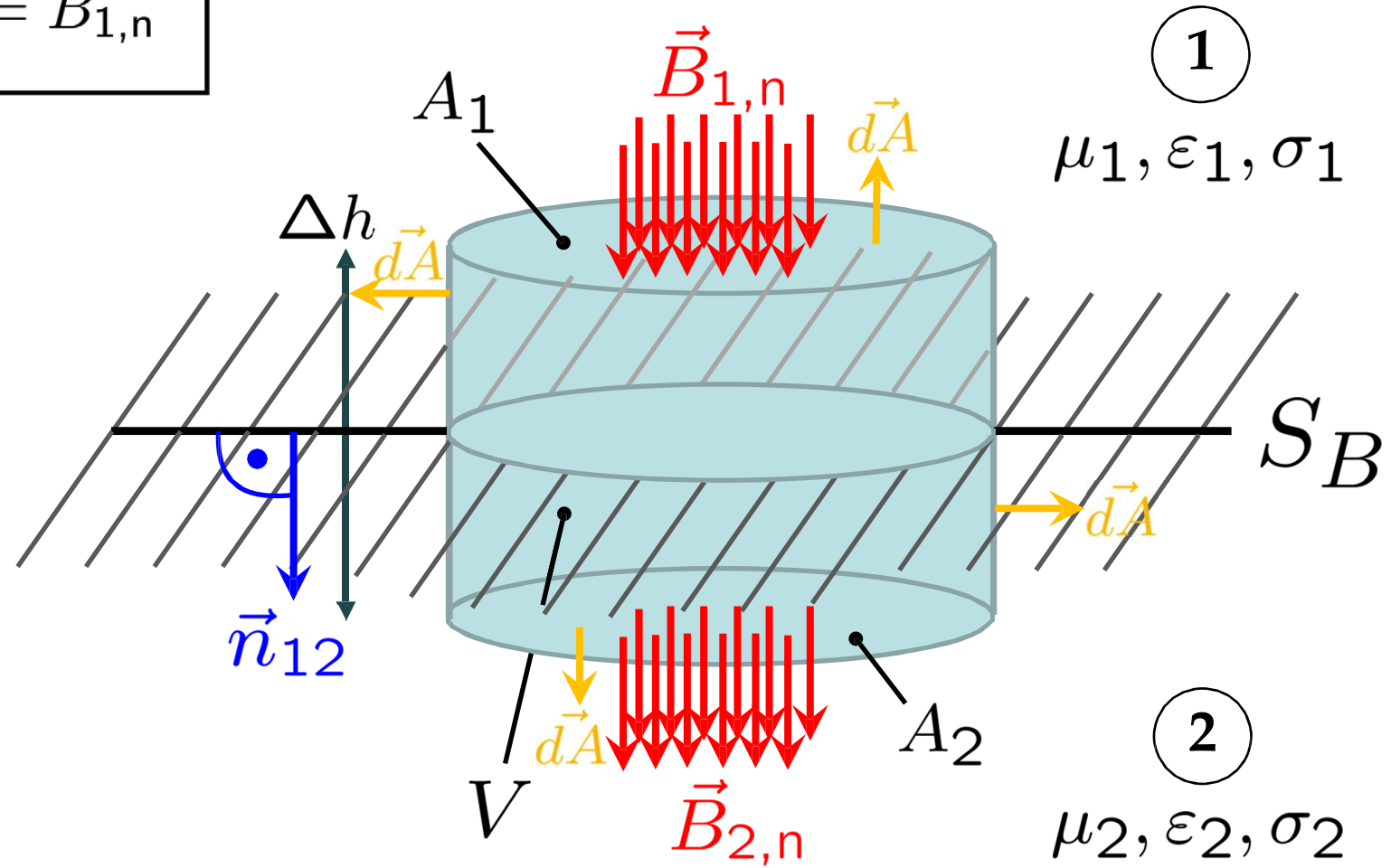
$$\vec{H}_{1,\text{tan}} = \vec{n}_{12} \times \vec{H}_1 \times \vec{n}_{12}$$

$$\vec{H}_{2,\text{tan}} = \vec{n}_{12} \times \vec{H}_2 \times \vec{n}_{12}$$

$$\vec{n}_{12} \times (\vec{H}_2 - \vec{H}_1) = \vec{J}_S$$

Boundary Conditions

$$B_{2,n} = B_{1,n}$$



PRACTICE EXERCISE 8.9

A unit normal vector from region 2 ($\mu = 2\mu_0$) to region 1 ($\mu = \mu_0$) is $\vec{a}_{n21} = (6\vec{a}_x + 2\vec{a}_y - 3\vec{a}_z)/7$. If $\vec{H}_1 = 10\vec{a}_x + \vec{a}_y + 12\vec{a}_z$ A/m and $\vec{H}_2 = H_{2x}\vec{a}_x - 5\vec{a}_y + 4\vec{a}_z$ A/m, determine

- H_{2x}
- The surface current density \vec{K} on the interface
- The angles θ_1 and θ_2 make with the normal to the interface

P.E. 8.9

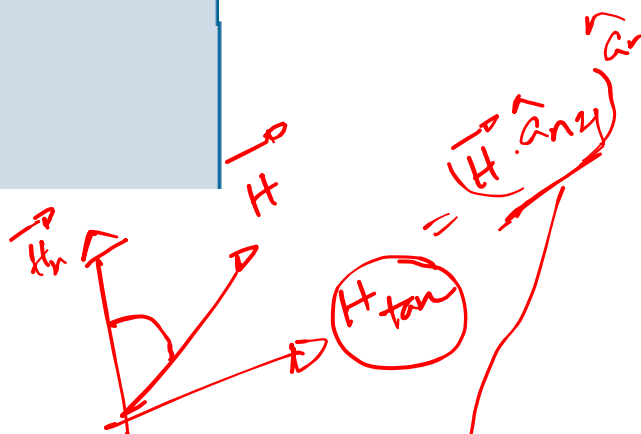
(a) $\vec{B}_{1n} = \vec{B}_{2n} \rightarrow \mu_1 \vec{H}_{1n} = \mu_2 \vec{H}_{2n}$

or $\mu_1 \vec{H}_1 \cdot \vec{a}_{n21} = \mu_2 \vec{H}_2 \cdot \vec{a}_{n21}$

$$\mu_0 \frac{(60 + 2 - 36)}{7} = 2\mu_0 \frac{(6H_{2x} + 10 - 12)}{7}$$

$$35 = 6H_{2x}$$

$$\underline{\underline{H_{2x} = 5.833}}$$



(b) $\vec{K} = (\vec{H}_1 - \vec{H}_2) \times \vec{a}_{n12} = \vec{a}_{n21} \times (\vec{H}_1 - \vec{H}_2)$

$$= \vec{a}_{n21} \times \left[(10, 1, 12) - \left(\frac{35}{6}, -5, 4 \right) \right]$$

$$= \frac{1}{7} \begin{vmatrix} 6 & 2 & -3 \\ 25/6 & 6 & 8 \end{vmatrix}$$

$$\underline{\underline{\vec{K} = 4.86\vec{a}_x - 8.64\vec{a}_y + 3.95\vec{a}_z \text{ A/m}}}$$

- (c) Since $\vec{B} = \mu\vec{H}$, \vec{B}_1 and \vec{H}_1 are parallel, i.e. they make the same angle with the normal to the interface.

$$\cos \theta_1 = \frac{\vec{H}_1 \cdot \vec{a}_{n21}}{|\vec{H}_1|} = \frac{26}{7\sqrt{100+1+144}} = 0.2373$$

$$\underline{\underline{\theta_1 = 76.27^\circ}}$$

$$\cos \theta_2 = \frac{\vec{H}_2 \cdot \vec{a}_{n21}}{|\vec{H}_2|} = \frac{13}{7\sqrt{(5.833)^2 + 25 + 16}} = 0.2144$$

$$\underline{\underline{\theta_2 = 77.62^\circ}}$$

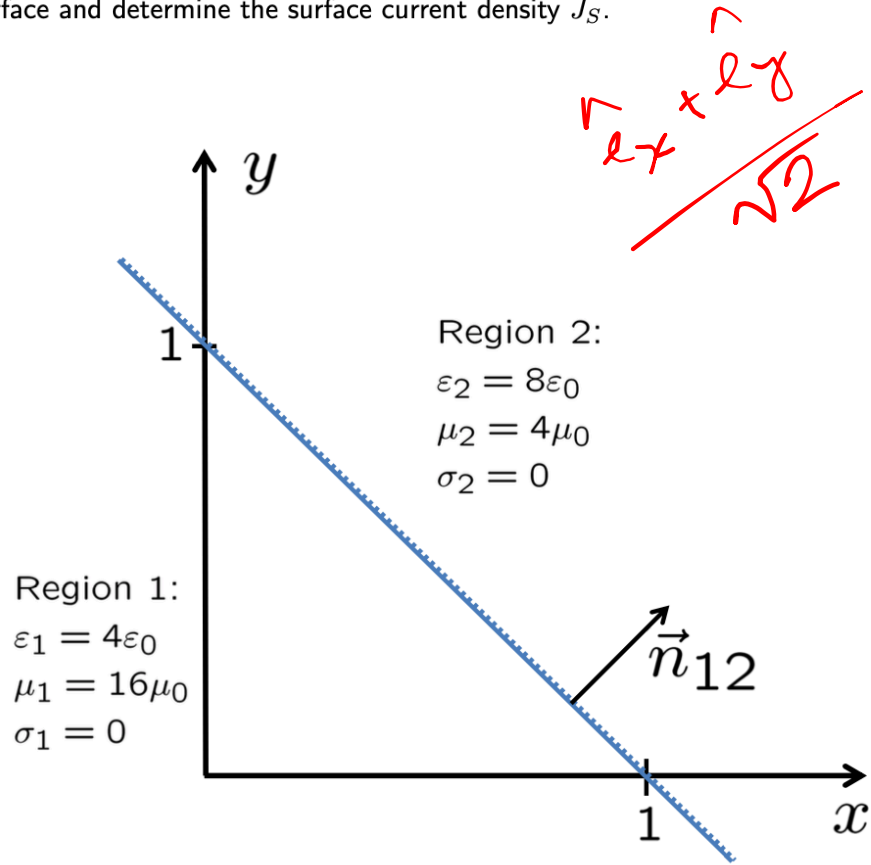
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- b) Find the magnetic induction \vec{B}_2 in region 2 so that the boundary conditions at the interface are fulfilled.
- c) Now the medium in region 2 is replaced by an ideal conductor with conductivity of $\sigma_2 \rightarrow \infty$ while all other parameters remain the same. Determine the surface charge density ρ_s on the interface and determine the surface current density \vec{J}_s .



3.4.6) $\vec{H}_{2tan} = \vec{H}_{1tan}$ because $\vec{J}_s = 0$

$\vec{H}_{1tan} = \vec{n}_{12} \times \vec{H}_1 \times \vec{n}_{12} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \times \begin{pmatrix} 0.1 \\ 0 \\ 0.2 \end{pmatrix} \times \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{A}{m}$

$= \frac{1}{\sqrt{2}} \begin{pmatrix} 0.2 \\ -0.2 \\ -0.1 \end{pmatrix} \times \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{A}{m} = \frac{1}{2} \begin{pmatrix} 0.1 \\ -0.1 \\ 0.4 \end{pmatrix} \frac{A}{m}$

$= \vec{H}_{2tan} \Rightarrow \vec{B}_{2tan} = \mu_2 \vec{H}_{2tan}$

$B_{2n} = B_{1n} \Rightarrow \mu_2 H_{2n} = \mu_1 H_{1n} \Rightarrow H_{2n} = \frac{\mu_1}{\mu_2} H_{1n} = 4 H_{1n}$

$H_{1n} = \vec{n}_{12} \cdot \vec{H}_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 0.1 \\ 0 \\ 0.2 \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{A}{m} = \frac{0.1}{\sqrt{2}} \frac{A}{m}$

$H_{2n} = 4 H_{1n} = \frac{0.4}{\sqrt{2}} \frac{A}{m} \Rightarrow \vec{H}_{2n} = \frac{0.4}{\sqrt{2}} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{A}{m} = \begin{pmatrix} 0.2 \\ 0.2 \\ 0 \end{pmatrix} \frac{A}{m}$

$\vec{B}_2 = \mu_2 \vec{H}_2 = \mu_2 (\vec{H}_{2tan} + \vec{H}_{2n})$

$= \mu_2 \left[\frac{1}{2} \begin{pmatrix} 0.1 \\ -0.1 \\ 0.4 \end{pmatrix} + \begin{pmatrix} 0.2 \\ 0.2 \\ 0 \end{pmatrix} \right] \frac{A}{m} = \mu_2 \begin{pmatrix} 0.25 \\ 0.15 \\ 0.2 \end{pmatrix} \frac{A}{m}$

3.3 Problem 3

The three dimensional space is split into two regions filled with different materials (see Fig. 3.3.2). The electric vector field \vec{E}_1 and the magnetic vector field \vec{H}_1 , in region 1 (see Fig. 3.3.2), are given as

$\vec{E}_1 = (2.0\vec{e}_y + 3.0\vec{e}_z) \text{ V/m}$ $\vec{H}_1 = (0.1\vec{e}_x + 0.2\vec{e}_z) \text{ A/m}$

For sub-problems a) and b) there are no surface charges ρ_s on the interface.

- a) Find the electric vector field \vec{E}_2 in region 2 so that the boundary conditions at the interface are fulfilled.
- b) Find the magnetic induction \vec{B}_2 in region 2 so that the boundary conditions at the interface are fulfilled.
- c) Now the medium in region 2 is replaced by an ideal conductor with conductivity of $\sigma_2 \rightarrow \infty$ while all other parameters remain the same. Determine the surface charge density ρ_s on the interface and determine the surface current density \vec{J}_s .

3.4c) $\vec{J}_s = \vec{n}_{12} \times (\vec{H}_2 - \vec{H}_1) = -\vec{n}_{12} \times \vec{H}_1 = \vec{H}_1 \times \vec{n}_{12}$

$= \begin{pmatrix} 0.1 \\ 0 \\ 0.2 \end{pmatrix} \times \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \frac{\text{A}}{\text{m}} = \begin{pmatrix} -\frac{0.2}{\sqrt{2}} \\ \frac{0.2}{\sqrt{2}} \\ \frac{0.1}{\sqrt{2}} \end{pmatrix} \frac{\text{A}}{\text{m}}$

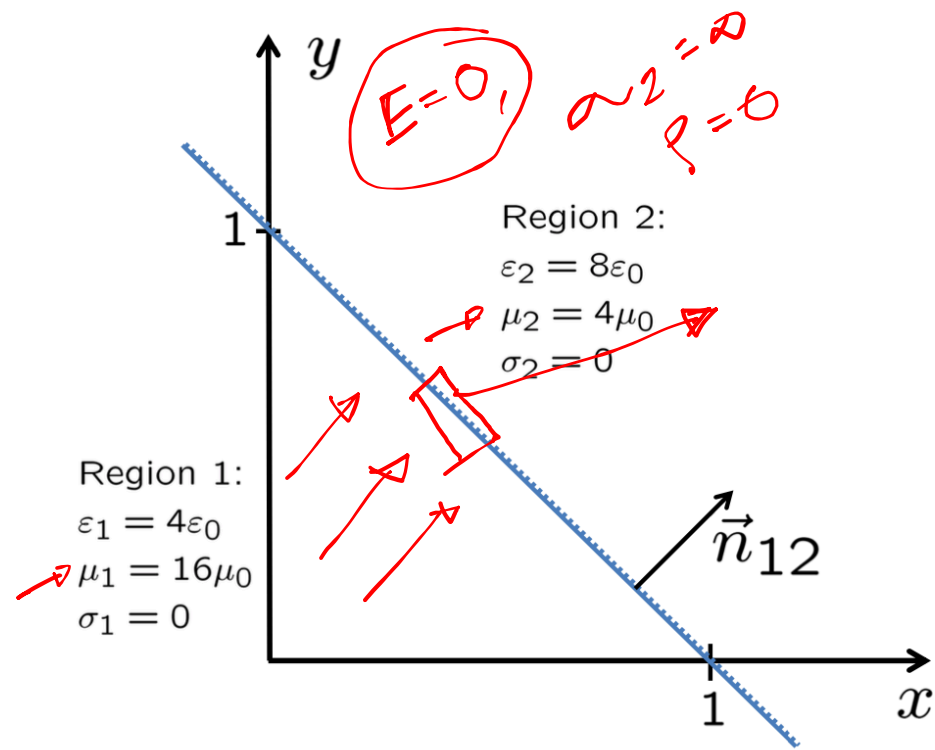
$\vec{D}_{1n} - \vec{D}_{2n} = \rho_s \Rightarrow \rho_s = -\vec{D}_{1n} = -\epsilon_1 E_{1n}$

$E_{1n} = \vec{n}_{12} \cdot \vec{E}_1 = \frac{1}{\sqrt{2}} \frac{\text{V}}{\text{m}}$

$\Rightarrow \rho_s = -\epsilon_1 \frac{1}{\sqrt{2}} \frac{\text{V}}{\text{m}}$ with $\epsilon_1 = 4\epsilon_0$

$= -4 \cdot 8.854 \cdot 10^{-12} \frac{\text{As}}{\text{Vm}}$

$\Rightarrow [\rho_s] = \frac{\text{As}}{\text{m}^2} = \frac{\text{C}}{\text{m}^2}$

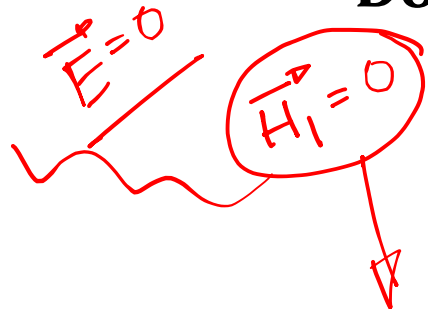


Maxwell's Equations

$$B_{2,n} = B_{1,n}$$

Boundary conditions at a perfect electric conductor:

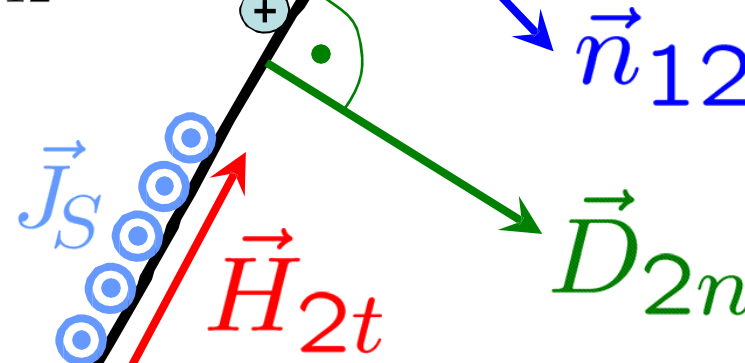
$$\vec{n}_{12} \times (\vec{H}_2 - \vec{H}_1) = \vec{J}_S$$



① $\mu_1, \epsilon_1, \sigma_1 \rightarrow \infty$ *cond*

$$\begin{aligned} \vec{n}_{12} \times \vec{H}_2 &= \vec{J}_S \\ \vec{H}_{2t} &= \vec{J}_S \times \vec{n}_{12} \\ |\vec{H}_{2t}| &= |\vec{J}_S| \end{aligned}$$

✓ $\vec{H}_{2,\text{tan}} = \vec{n}_{12} \times \vec{H}_2 \times \vec{n}_{12}$



② $\mu_2, \epsilon_2, \sigma_2$

$B_{2n} = 0$

Maxwell's Equations

Boundary conditions at a perfect electric conductor:

$$\sigma_1 \rightarrow \infty$$

$$\vec{n}_{12} \times \vec{H}_2 = \vec{J}_S$$

$$\vec{H}_{2t} = \vec{J}_S \times \vec{n}_{12}$$

$$|\vec{H}_{2t}| = |\vec{J}_S|$$

$$\vec{E}_{2t} = \vec{0}$$

$$D_{2n} = \rho_S$$

$$B_{2n} = 0$$

3.5 Problem 5

Two very large metallic plates with an area A are placed next to each other in a distance d (see Fig. 3.5.4).

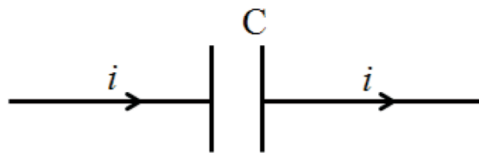


Figure 3.5.4

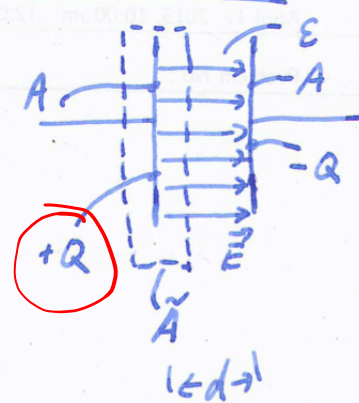
Derive the capacitance C between the plates.

Hint: Assume that one plate carries positive and the other plate the same amount of negative charges. Additionally assume a pure homogeneous field (because the plates are very large).

$$C = \frac{\epsilon A}{d}$$

$$V = E \times d$$

Problem 3.5



$$Q = C V \Rightarrow C = \frac{Q}{V} \quad (12)$$

$$Q = ? ; V = ?$$

$$\oint \vec{D} \cdot d\vec{A} = \iiint \rho dV = Q$$

here: $\vec{D} \parallel d\vec{A}$ and \vec{D} is constant (homogeneous) between the plates and zero elsewhere.

$$\Rightarrow |\vec{D}| \cdot A = Q ; |\vec{D}| = \epsilon |\vec{E}|$$

$$\Rightarrow Q = \epsilon |\vec{E}| \cdot A$$

$$V = \int_{\text{plate 1}}^{\text{plate 2}} \vec{E} \cdot d\vec{s} ; \text{ here: } \vec{E} \parallel d\vec{s} \text{ and } \vec{E} \text{ is constant (homogeneous) between the plates.}$$

$$\Rightarrow |\vec{E}| \cdot d = V$$

$$\Rightarrow \text{Capacitance } C = \frac{Q}{V} = \frac{\epsilon |\vec{E}| A}{|\vec{E}| d}$$

$$= \epsilon \frac{A}{d} \checkmark$$